# On Riemannian manifolds endowed with a $\mathcal{T}$-parallel almost contact 4 -structure 

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#### Abstract

T}\)-parallel almost contact 4-structures on a Riemannian manifold are studied. It is proved that such a manifold is a local Riemannian product of two totally geodesic submanifolds, one of them being a space form. Additional results are obtained when the manifold is endowed with a framed $f$-structure.


## 1. Introduction

In the last two decades, contact, almost contact, paracontact and almost cosymplectic manifolds carrying $r(r>1)$ Reeb vector fields $\xi_{r}$ have been studied by a certain number of authors, as for instance: M. Kobayashi [11], A. Bucki [4], S. Tachibana and W. N. Yu [22], K. Yano and M. Kon [25], V. V. Goldberg and R. Rosca [8] and some others.

In the present paper we consider a $(2 m+4)$ dimensional Riemannian manifold carrying 4 structure vector fields $\xi_{r}(r, s \in\{2 m+1, \ldots, 2 m+4\})$ and with a distinguished vector field $\mathcal{T}$, such that the vertical connection forms define a $\mathcal{T}$-parallel connection and the Reeb vector fields are $\mathcal{T}$-parallel (this structure is called a $\mathcal{T}$-parallel almost contact 4 -structure and it will be defined in Definition 3.1). Then we shall prove that such a manifold is a local Riemannian product of two totally geodesic submanifolds, $M=M^{\top} \times M^{\perp}$, where $M^{\perp}$ is a space form tangent to the distribution generated by the Reeb vector fields, and that the vector field $\mathcal{T}$ is closed torse forming (Theorem 3.3).

In section 4 we shall study conformal-type structures induced by a $\mathcal{T}$ parallel almost contact 4 -structure. Finally, in section 5 we assume that the manifold under consideration is endowed with a framed $f$-structure, proving that $M^{\top}$ is a Kählerian submanifold (Theorem 5.2).

## 2. Preliminaries

Let $(M, g)$ be a Riemannian $C^{\infty}$-manifold and let $\nabla$ be the covariant differential operator defined by the metric tensor $g$. We assume that $M$ is oriented and $\nabla$ is the Levi-Civita connection. Let $\Gamma T M$ be the set of sections of the tangent bundle $T M$ and $b: T M \rightarrow T^{*} M, X \rightarrow X^{b}$, the musical isomorphism defined by $g$. Next, following a standard notation, we set: $A^{q}(M, T M)=\operatorname{Hom}\left(\Lambda^{q} T M, T M\right)$ and notice that elements of $A^{q}(M, T M)$ are vector valued $q$-forms $(q \leq \operatorname{dim} M)$. Denote by $d^{\nabla}$ : $A^{q}(M, T M) \rightarrow A^{q+1}(M, T M)$ the exterior covariant derivative operator with respect to $\nabla$ (it should be noticed that generally $d^{\nabla^{2}}=d^{\nabla} \circ d^{\nabla} \neq 0$, unlike $\left.d^{2}=d \circ d=0\right)$. The identity tensor field $I$ of type $(1,1)$ can be considered as a vector valued 1 -form $I \in A^{1}(M, T M)$ (and it is also called the soldering form [7]).

We shall remember the following
Definition 2.1. (1) (see [10]) The operator $d^{\omega}=d+e(\omega)$ acting on $\Lambda M$ is called the cohomology operator, where $e(\omega)$ means the exterior product by the closed 1-form $\omega \in \Lambda^{1} M$, i.e., $d^{\omega} u=d u+\omega \wedge u$ for any $u \in \Lambda M$. One has $d^{\omega} \circ d^{\omega}=0$, and if $d^{\omega} u=0, u$ is said to be $d^{\omega}$-closed. If $\omega$ is exact, then $u$ is said to be $d^{\omega}$-exact.
(2) (see [18], [16]) Any vector field $X \in \Gamma T M$ such that: $d^{\nabla}(\nabla X)=$ $\nabla^{2} X=\pi \wedge I \in A^{2}(M ; T M)$ for some 1-form $\pi$, is called an exterior concurrent vector field and the 1 -form $\pi$, which is called the concurrence form, given by $\pi=f X^{b}, f \in C^{\infty}(M)$.
(3) (see [23], [16]) A vector field $\mathcal{T}$ whose covariant differential satisfies $\nabla \mathcal{T}=r I+\alpha \otimes \mathcal{T} ; r \in C^{\infty}(M)$ where $\omega=\mathcal{T}^{b}$ is a closed form, is called a closed torse forming.

If $\Re$ denotes the Ricci tensor of $\nabla$ and $X$ an exterior concurrent vector field, one has $\Re(X, Z)=-(n-1) f g(X, Z), Z \in \Gamma T M, n=\operatorname{dim} M$.

Let $C$ be any conformal vector field on $M$ (i.e., the conformal version of Killing's equations). As is well known, $C$ satisfies

$$
\begin{equation*}
\mathcal{L}_{C} g(C, Z)=\rho g(C, Z) \text { or } g\left(\nabla_{Z} C, Z^{\prime}\right)+g\left(\nabla_{Z} C, Z\right)=\rho g\left(Z, Z^{\prime}\right) \tag{2.1}
\end{equation*}
$$

$\left(Z, Z^{\prime} \in \Gamma T M\right)$ where the conformal scalar $\rho$ is defined by $\rho=\frac{2}{n}(\operatorname{div} C)$.
We recall the following basic formulas (see [3])

Proposition 2.2. With the above notation, let $\mathcal{L}_{C}, K, \Delta$ and $\Re$ denote the Lie derivative with respect to $C$, the scalar curvature, the Laplacian and the Ricci tensor field of $\nabla$, respectively. Then:
(1) $\mathcal{L}_{C} Z^{b}=\rho Z^{b}+[C, Z]^{b}$ (Orsted's lemma).
(2) $\mathcal{L}_{C} K=(n-1) \Delta \rho-K \rho$.
(3) $2 \mathcal{L}_{C} \Re\left(Z, Z^{\prime}\right)=\Delta \rho g\left(Z, Z^{\prime}\right)-(n-2)\left(\operatorname{Hess}_{\nabla} \rho\right)\left(Z, Z^{\prime}\right)$ where $\left(\operatorname{Hess}_{\nabla} \rho\right)\left(Z, Z^{\prime}\right)=g\left(Z, \nabla_{Z^{\prime}}(\operatorname{grad} \rho)\right)$.
Definition 2.3 (see [19], [20], [15]). Any vector field $C$ whose covariant differential satisfies $\nabla C=f I+C \wedge X$ is said to be a skew-symmetric conformal (ab. SKC) vector field or a structure conformal vector field, where $\wedge$ means the wedge product of vector fields, i.e., $(X \wedge Y) Z=$ $g(Y, Z) X-g(X, Z) Y ; X, Y, Z \in \Gamma(T M)$.

Remark 2.4. Let $\mathcal{O}=\operatorname{vect}\left\{e_{A} ; A \in 1, \ldots, n\right\}$ be an adapted local field of orthonormal frames on $M$ and let $\mathcal{O}^{*}=\operatorname{covect}\left\{\omega^{A}\right\}$ be the associated coframe. With respect to $\mathcal{O}$ and $\mathcal{O}^{*}$ the soldering form $I$ and E. Cartan's structure equations can be written in indexless manner as
(1) $I=\omega^{A} \otimes e_{A} \in A^{1}(M, T M)$
(2) $\nabla e=\vartheta \otimes e \in A^{1}(M, T M)$
(3) $d \omega=-\vartheta \wedge \omega$
(4) $d \vartheta=-\vartheta \wedge \vartheta+\Theta$

In the above equations $\vartheta$ (resp. $\Theta$ ) are the local connection forms in the bundle $O(M)$ (resp. the curvature form on $M$ ).

Finally, we remember the following
Proposition 2.5. Let $\pi \in \Lambda^{1} M$ be a Pffaf form on a manifold $M$. Then in order that $\pi$ be of class $2 s$ on $M$ it is necessary and sufficient to have $(d \pi)^{s+1}=0, \pi \wedge(d \pi)^{s}=0$.

## 3. The main result

Let $M\left(\xi_{r}, \eta^{r}, g\right)$ be a $(2 m+4)$-dimensional oriented Riemannian manifold carrying 4 Reeb vector fields $\xi_{r}(r, s \in\{2 m+1, \ldots, 2 m+4\})$ with associated structure covectors $\eta^{r}$, that is $\eta^{r}\left(\xi_{s}\right)=\delta_{r s}$. Following a known terminology we may decompose the tangent space $T_{p}(M)$ at $p \in M$ to $M$ as $T_{p} M=D_{p}^{\top} \oplus D_{p}^{\perp}$. Then $D_{p}^{\perp}$ is a 4-dimensional distribution defined by the set $\left\{\xi_{r}\right\}$, called the vertical distribution, and its orthogonal complement $D_{p}^{\top}=\left\{\xi_{r}\right\}^{\perp}$ which is called the horizontal distribution. Consequently any vector field $Z \in \Gamma(T M)$ may be written as $Z=\left(Z-\eta^{r}(Z) \xi_{r}\right)+\eta^{r}(Z) \xi_{r}=Z^{\top}+Z^{\perp}$ where $Z^{\top}$ (resp. $Z^{\perp}$ ) is the
horizontal component of $Z$ (resp. the vertical component of $Z$ ). We recall that setting $A ; B \in\{1,2, \ldots, 2 m\}$ the connection forms $\vartheta_{B}^{A}, \vartheta_{B}^{r}$ and $\vartheta_{s}^{r}$ are called the horizontal, the transversal and the vertical connection forms respectively (see also [21]).

With the above notation, one has the following
Definition 3.1 ([17], [9]). Let $M\left(\xi_{r}, \eta^{r}, g\right)$ be a $(2 m+4)$-dimensional oriented Riemannian manifold carrying 4 Reeb vector fields $\xi_{r}$ such that the vertical connection forms verifies $\vartheta_{s}^{r}=\left\langle\mathcal{T}, \xi_{s} \wedge \xi_{r}\right\rangle$, where $\mathcal{T}$ is a certain vertical vector field. Then, we say that vertical connection forms $\vartheta_{s}^{r}$ define on $D^{\perp}$ a $\mathcal{T}$-parallel connection and $\mathcal{T}$ is called the generator of the considered ( $\mathcal{T} . P)$-connection. Moreover, if the Reeb vector fields are $\mathcal{T}$-parallel, i.e., $\nabla_{\mathcal{T}} \xi_{r}=0$, then the manifold $M\left(\xi_{r}, \eta^{r}, g\right)$ is said to be endowed with a $\mathcal{T}$-parallel almost contact 4 -structure (abr. $\mathcal{T} . P . A . C .4$-structure).

In the present paper we shall deal with these manifolds.
Remark 3.2. If we set $\mathcal{T}=\sum t_{r} \xi_{r} ; t_{r} \in C^{\infty}(M)$ then the vertical connection forms are expressed by $\vartheta_{s}^{r}=t_{s} \eta^{r}-t_{r} \eta^{s}$. Since the vertical connection forms satisfy $\vartheta_{s}^{r}(\mathcal{T})=0$, then by reference to [13] we may say that $\vartheta_{s}^{r}$ are relations of integral invariance for the vector field $\mathcal{T}$.

Similarly one may decompose in an unique fashion the soldering form $I$ of $M$ as $I=I^{\top}+I^{\perp}$ where $I^{\top}=\omega^{A} \otimes e_{A}$ and $I^{\perp}=\eta^{r} \otimes \xi_{r}$ mean the line element of $D^{\top}$ and the line element of $D^{\perp}$ respectively.

We can state
Theorem 3.3. Let $M\left(\xi_{r}, \eta^{r}, g\right)$ be a $(2 m+4)$-dimensional Riemannian manifold endowed with a $\mathcal{T}$-parallel almost contact 4 -structure and let $\mathcal{T}$ be the generator vector field of this structure.

For such a manifold the structure covectors $\eta^{r}(r \in\{2 m+1, \ldots, 2 m+4\})$ are of class 2 and cohomologically exact, i.e., $d^{-\omega} \eta^{r}=0$, where $\omega$ is the dual form of the generator $\mathcal{T}$ which enjoys the property to be a closed torse forming and to define a relative infinitesimal conformal transformation of the almost contact structure of $M$.

Any manifold $M$ which carries a (T .P.A.C.) 4-structure may be viewed as the local Riemannian product $M=M^{\top} \times M^{\perp}$ such that:
(i) $M^{\perp}$ is a totally geodesic submanifold of $M$, tangent to the vertical distribution $D^{\perp}=\left\{\xi_{r}\right\}$ which enjoys the property to be a space form of curvature $-2 a$ ( $a=$ const.)
(ii) $M^{\top}$ is a totally geodesic submanifold of $M$, tangent to the horizontal distribution $D^{\top}=\left\{\xi_{r}\right\}^{\perp}$ of $M$.

Proof. Making use of the structure equations of Remark 2.4(2) and taking account of Remark 3.2 one derives:

$$
\begin{equation*}
\nabla \xi_{r}=t_{r} I^{\perp}-\eta^{r} \otimes \mathcal{T} \tag{3.1}
\end{equation*}
$$

Hence if $Z_{1}^{\perp}, Z_{2}^{\perp} \in D_{p}^{\perp}$ are any vertical vector fields, it quickly follows from (3.1) $\nabla_{Z_{2}^{\perp}} Z_{1}^{\perp} \in D_{p}^{\perp}$. This, as is known, proves that $D_{p}^{\perp}$ is an autoparallel foliation and that the leaves $M^{\perp}$ of $D_{p}^{\perp}$ are totally geodesic submanifolds of $M$ (in our case, $\operatorname{dim} M^{\perp}=4$ ). Next making use of the structure equations of Remark 2.4(3) one finds

$$
\begin{equation*}
d \eta^{r}=\omega \wedge \eta^{r} \tag{3.2}
\end{equation*}
$$

where $\omega=\mathcal{T}^{b}$ denotes the dual form of the generator vector field $\mathcal{T}$.
By reference to [7], equations (3.2) show that all the Reeb covectors $\eta^{r}$ are exterior recurrent and by a simple argument it follows that the recurrence form $\omega$ is necessarly closed, i.e., $d \omega=0$. With the help of (3.1) and (3.2) one also derives from $I^{\perp}=\eta^{r} \otimes \xi_{r}$ that $I^{\perp}$ is exterior covariant closed, i.e., $d^{\nabla}\left(I^{\perp}\right)=0$ and this is matching the fact that $I^{\perp}$ is the soldering form of the leaf $M^{\perp}$. By reference to Proposition 2.5 it is seen by (3.2) that the structure covectors $\eta^{r}$ are of class 2 .

Let now denote by $\varphi=\eta^{2 m+1} \wedge \ldots \wedge \eta^{2 m+4}$ the simple form which coresponds to $D_{p}^{\perp}$ (or equivalently the volume element of $M^{\perp}$ ). By (3.2) one has at once $d \varphi=0$ and therefore since one may write $D_{p}^{\top} \subset \operatorname{ker}(\varphi) \cap$ $\operatorname{ker}(d \varphi)$ we conclude that the horizontal distribution $D_{p}^{\top}$ is also involutive. Then setting $M^{\top}$ for the $2 m$ leaf of $D_{p}^{\top}$, it is seen that $\xi_{r}$ are geodesic normal section for the immersion $\kappa: M^{\top} \rightarrow M$, which is totally geodesic. It follows from the above discussion that the manifold $M$ under consideration is the local product $M=M^{\top} \times M^{\perp}$, where $M^{\top}$ and $M^{\perp}$ are totally geodesic submanifolds of $M$, tangent to the horizontal distribution $D^{\top}$ and the vertical distribution $D^{\perp}$ of $M$ respectively.

Further since the dual form $\omega$ of $\mathcal{T}$ is expressed by $\omega=t_{r} \eta^{r}$ then by virtue of (3.2) one may set

$$
\begin{equation*}
d t_{r}=\lambda \eta^{r} \Longrightarrow d \lambda-\lambda \omega=0 \tag{3.3}
\end{equation*}
$$

which shows that $\omega$ is an exact form. In consequence of this fact, equations (3.2) may be expressed, using the notation introduced in Definition 2.1(1),
as $d^{-\omega} \eta^{r}=0$, thus proving that the structure covectors of $M\left(\xi_{r}, \eta^{r}, g\right)$ are cohomologically exact.

Taking now the covariant differential of the generator vector field $\mathcal{T}$, one derives on behalf of (3.1) and (3.3)

$$
\begin{equation*}
\nabla \mathcal{T}=(\lambda+2 t) I^{\perp}-\nu \otimes \mathcal{T} ; 2 t=\|\mathcal{T}\|^{2} \tag{3.4}
\end{equation*}
$$

which shows the significative fact that $\mathcal{T}$ is a closed torse forming (def. 2.1(3)). Since this quality implies that $\mathcal{T}$ is a gradient vector field, this fact is in accordance with equation (3.3). We also derive from (3.4)

$$
\begin{equation*}
d t=\lambda \omega \Longrightarrow t+\lambda=a=\text { const. } \tag{3.5}
\end{equation*}
$$

Next operating on (3.1) by the exterior covariant derivative operator $d^{\nabla}$ one quickly derives by (3.2) and (3.4) that one has $d^{\nabla}\left(\nabla \xi_{r}\right)=\nabla^{2} \xi_{r}=2 a \eta^{r} \wedge$ $I^{\perp}$. The above equations reveal the important fact that all the vectors $\left\{\xi_{r}\right\}$ on $M^{\perp}$ are exterior concurrent vector fields (see [20]). Then since the conformal scalar $2 a$ is constant, we conclude by reference to [16] that the vertical submanifold $M^{\perp}$ is a space form of curvature $-2 a$.

Next by (3.2), (3.3) and (3.5) one derives succesively $\mathcal{L}_{\mathcal{T}} \eta^{r}=$ $(a+t) \eta^{r}-t_{r} \omega$ and $d\left(\mathcal{L}_{\mathcal{T}} \eta^{r}\right)=(2 a+\lambda) \omega \wedge \eta^{r}$. In consequence of the last equation and by reference to [14] we agree to say that the generator vector $\mathcal{T}$ defines a relative infinitesimal conformal transformation of the considered almost contact 4 -structure, thus finishing the proof.

## 4. Conformal-type structures induced by a (T.P.A.C.) 4 -structure

In the present section we consider on $M^{\perp}$ the 2 -form $\psi$ of rank 2 (if $\Omega \in \Lambda^{2} M, \operatorname{rank} r$ is the smallest integer such that $\Omega^{r+1}=0$ ), defined by $\psi=\eta^{2 m+1} \wedge \eta^{2 m+2}+\eta^{2 m+3} \wedge \eta^{2 m+4}$. On behalf of (3.2) one quickly derives by exterior differentiation of $\psi$ that $d \psi=2 \omega \wedge \psi \Leftrightarrow d^{-2 \omega} \psi=0$ (the last equality obtained on behalf of Definition 2.1(1)). Therefore following a known definition it is seen that $\psi$ is a conformal symplectic form on $M^{\perp}$ having $\omega$ (resp. $\mathcal{T}$ ) as covector of Lee (resp. vector field of Lee). In addition in the case under discussion one may say that $\psi$ is a $d^{-2 \omega}$-exact form.

It should be noticed that this property is in accordance with the general properties of $\mathcal{T}$-parallel connections (see also [14]). If $Y \in \Gamma T M^{\perp}$ is any vertical vector field, then by reference to [12] we set ${ }^{b} Y=-i_{Y} \psi$. Do not confuse with the the musical isomorphism $b: \Gamma T M \rightarrow \Gamma T M^{*}$, which is denoted by $X \rightarrow X^{b}$. For instance, $\omega=\mathcal{T}^{b}$.

In the case under discussion and in order to simplify we write

$$
\beta=-{ }^{b} \mathcal{T}=t_{2 m+1} \eta^{2 m+2}+t_{2 m+3} \eta^{2 m+4}-t_{2 m+2} \eta^{2 m+1}-t_{2 m+4} \eta^{2 m+3}
$$

and by (3.3) and (3.2) one gets $d \beta=2 \lambda \psi+\omega \wedge \beta$ by which after a standard calculation one derives $\mathcal{L}_{\mathcal{T}} \psi=2(a+t) \psi-\omega \wedge \beta$. Since $\omega$ is an exact form, then following [1] the above equation shows that $\mathcal{T}$ defines a weak infinitesimal conformal transformation of $\psi$. Then we obtain $d\left(\mathcal{L}_{\mathcal{T}} \psi\right)=$ $8 a \omega \wedge \psi$. Therefore we may also say that $\mathcal{T}$ defines a relative infinitesimal conformal transformation of $\psi$.

Consider now the vertical vector field $C=C^{r} \xi_{r}$ and set $\varrho={ }^{b} C$. Then in order that $C$ be an infinitesimal conformal transformation of $\psi$, one finds making use of (3.2)

$$
\begin{equation*}
d C^{r}=C^{r} \omega \tag{4.1}
\end{equation*}
$$

This implies $d \varrho=2 \omega \wedge \varrho \Leftrightarrow d^{-2 \omega} \varrho=0$ and setting $s=g(C, \mathcal{T})$ one may write $\mathcal{L}_{\mathcal{T}} \psi=2 s \psi$. In the light of this problem, and making use of (3.1) and (4.1) one derives

$$
\begin{equation*}
\nabla C=s I^{\perp}+C \wedge \mathcal{T} \tag{4.2}
\end{equation*}
$$

which reveals the important fact that $C$ is a structure conformal vector field having $2 s=\rho$ as conformal scalar (see Definition 2.3). Setting $\alpha=C^{b}$ one finds by (3.4) and (4.2)

$$
\begin{equation*}
d s=\lambda \alpha+s \omega \tag{4.3}
\end{equation*}
$$

and on the other hand by (3.2) one has

$$
\begin{equation*}
d \alpha=2 \omega \wedge \alpha \Longleftrightarrow d^{-2 \omega} \alpha=0 \tag{4.4}
\end{equation*}
$$

Hence one may say that as $\psi$ the dual form $\alpha$ of $C$ is $d^{-2 \omega}$-exact. It should be noticed that equation (4.4) is in accordance with the general properties of structure conformal vector fields [19] (see also [14], [15]).

By (3.3), (4.3) and (4.4) it is seen that the existence of the structure conformal vector field $C$ is determined by the exterior differential system $\Sigma_{e}$ whose characterisitic numbers are $r=3, s_{0}=2, s_{1}=1$. Since $r=$ $s_{0}+s_{1}$ it follows by E. Cartan's test [5] that $\Sigma_{e}$ is involutive and $C$ is determined by 1 arbitrary function of 2 arguments.

Next since $\rho=2 s$, it follows at once from (4.3), by duality: $\operatorname{grad} \rho=$ $2 \lambda C+\rho \mathcal{T}$. But as it is known $\operatorname{div} Z=\operatorname{tr}[\nabla Z], Z \in \Gamma T M$, and so one gets from (3.4) $\operatorname{div} \mathcal{T}=4 a+2 t$ and $C$ being a conformal vector field one has
$\operatorname{div} C=4 \rho$. Therefore by the general formulas $\Delta f=-\operatorname{div}(\operatorname{grad} f), f \in$ $C^{\infty} M$, a short calculation gives

$$
\begin{equation*}
\Delta \rho=-8 a \rho \tag{4.5}
\end{equation*}
$$

which shows that $\rho$ is an eigenfunction of $\Delta$ and has $-8 a$ associated eigen value. Following a known theorem, it follows that if $M^{\perp}$ is compact, then necessarily $a=-\mu^{2}$ ( $\mu=$ const.), that is, $M^{\perp}$ is an elliptic submanifold of $M$.

On the other hand taking the covariant differential of $\operatorname{grad} \rho$, then by a standard calculation one infers

$$
\begin{equation*}
\nabla \operatorname{grad} \rho=4 a \rho I^{\perp} \tag{4.6}
\end{equation*}
$$

which reveals that $\operatorname{grad} \rho$ is concurrent vector field on $M^{\perp}[6]$ (we recall that concurrency is of conformal nature). Accordingly on behalf of the definition given in [14], we may say in the case under consideration $C$ has the divergence conformal property. It is worth to point out that if $M^{\perp}$ is an elliptic submanifold of $M$ (i.e., $a=-\mu^{2}$ ), then following Obata's theorem [24], $M^{\perp}$ is isometric to a sphere of radius $\frac{1}{2} \mu$.

Further since $M^{\perp}$ is a space form, then we recall [16] that any vector field on $M^{\perp}$ is E.C., with the same conformal scalar $2 a$. Consequently, if $\Re$ denotes the Ricci tensor of $\nabla$, one has

$$
\begin{equation*}
\Re(C, Z)=-6 a g(C, Z), Z \in \Gamma T M^{\perp} \tag{4.7}
\end{equation*}
$$

Then by (4.5), (4.6), (4.7) and making use of Proposition 2.2(3) and carrying out the calculations one derives $\mathcal{L}_{C} g(C, Z)=\frac{4}{3} \rho g(C, Z)$. Therefore one may state that the (S.C)-vector field $C$ defines an infinitesimal conformal transformation of all the functions $g(C, Z)$ where $Z \in \Gamma T M^{\perp}$. It should be noticed that this situation is similar to that of [14]. In addition by (3.1) and (4.2) one finds

$$
\begin{equation*}
\left[C, \xi_{r}\right]=-\frac{\rho}{2} \xi_{r} \tag{4.8}
\end{equation*}
$$

which shows that the structure vector fields $\xi_{r}$ admit infinitesimal transformations of generator $C$. Next making use of Orsted's lemma (Proposition 2.2(1)) it follows

$$
\begin{equation*}
\mathcal{L}_{C} \eta^{r}=\rho \eta^{r} . \tag{4.9}
\end{equation*}
$$

Hence making use of a known terminology, it follows that $C$ defines an almost contact transformation of the structure covectors $\eta^{r}$.

Finally we denote by $\mathcal{P}=\xi_{2 m+1} \wedge \xi_{2 m+2}+\xi_{2 m+3} \wedge \xi_{2 m+4}$ the Poisson bivector [12] associated with the conformal symplectic form $\psi$. Since $\mathcal{P}$ may be expressed as

$$
\begin{aligned}
\mathcal{P}= & \eta^{2 m+2} \otimes \xi_{2 m+1}-\eta^{2 m+1} \otimes \xi_{2 m+2} \\
& +\eta^{2 m+4} \otimes \xi_{2 m+3}-\eta^{2 m+3} \otimes \xi_{2 m+4}
\end{aligned}
$$

then since the Lie derivative is additive, one gets by (4.8) and (4.9) that $\mathcal{L}_{C} \mathcal{P}=0$ which shows that $C$ defines an infinitesimal automorphism of $\mathcal{P}$.

Next operating on the vector valued 1 -form $\mathcal{P}$ by the operator $d^{\nabla}$ one derives after two sucesive computations $d^{\nabla} \mathcal{P}=\omega \wedge \mathcal{P}-2 \psi \otimes \mathcal{T}-$ $\beta \wedge I^{\perp} \in A^{2}(M, T M)\left(\beta=-{ }^{b} \mathcal{T}\right)$ and $d^{\nabla^{2}} \mathcal{P}=4 a \psi \wedge I^{\perp}$. Therfore (see Proposition 2.5) the last equality shows that $\mathcal{P}$ is a 2 -exterior vector valued 1-form. Moreover, taking into account $\mathcal{L}_{\mathcal{T}} \psi=2 s \psi$ a short calculation gives $\mathcal{L}_{C}\left(d^{\nabla^{2}} \mathcal{P}\right)=\frac{\rho}{2} d^{\nabla^{2}} \mathcal{P}$ that is $C$ defines an infinitesimal conformal transformation of $d^{\nabla^{2}} \mathcal{P}$.

Then one has the
Theorem 4.1. Let $M\left(\xi_{r}, \eta^{r}, g\right)$ be a $(2 m+4)$-dimensional Riemannian manifold endowed with a ( $\mathcal{T}$.P.A.C.) 4-structure discussed in Section 2 and having $\mathcal{T}$ as generator vector field. Let $M^{\perp}$ be the space form submanifold of $M$, tangent to the vertical distribution $D^{\perp}=\left\{\xi_{r}\right\}$ of $M$. One has the following properties:
(i) $M^{\perp}$ is equipped with a conformal symplectic structure $\operatorname{CSp}(4, \mathbf{R})$ defined by the form $\psi \in \Lambda^{2} M^{\perp}$ (of rank 2) and such that the covector of Lee corresponding to $\operatorname{CSp}(4, \mathbf{R})$ is the dual form $\omega$ of $\mathcal{T}$, that is, $d \psi=2 \omega \wedge \psi$ and $\mathcal{T}$ defines a relative infinitesimal conformal transformation of $\psi$, that is, $d\left(\mathcal{L}_{\mathcal{T}} \psi\right)=8 a \omega \wedge \psi,(a=$ const. $)$
(ii) Any vector field $C$ which defines an infinitesimal conformal transformation of $\psi$ is a structure conformal vector field, i.e., $\nabla C=g(\mathcal{T}, C) I^{\perp}+$ $C \wedge \mathcal{T}$ and one has $\mathcal{L}_{C} \psi=\rho \psi ; \rho=2 g(\mathcal{T}, C)$ and $\mathcal{L}_{C} g(C, Z)=\frac{4}{3} \rho g(C, Z)$, $Z \in \Gamma T M^{\perp}$.
(iii) The conformal scalar $\rho\left(\mathcal{L}_{C} g=\rho g\right)$ is an eigenfunction of $\Delta$ and if $M^{\perp}$ is compact, then $a=-\mu^{2}$ and $M^{\perp}$ is isometric to a sphere of radius $\frac{1}{2} \mu$.
(iv) The Poisson bivector $\mathcal{P}$ associated with $\psi$ is a 2-exterior vector valued 1-form, i.e., $d^{\nabla^{2}} \mathcal{P}=4 a \psi \wedge I^{\perp}$ and $C$ defines an infinitesimal automorphism of $\mathcal{P}$.

## 5. Framed $f$-structures

In the present section we assume that the manifold $M\left(\xi_{r}, \eta^{r}, g\right)$ under consideration is endowed with a framed $f$-structure $\phi[27]$, that is $\phi$ is a tensor field of type $(1,1)$ and rank $2 m$ which satisfies:
(1) $\phi^{3}+\phi=0$
(2) $\phi^{2}=-I+\sum \eta^{r} \otimes \xi_{r} ; \phi \xi_{r}=0 ; \eta^{r} \circ \phi=0$
(3) $g\left(Z, Z^{\prime}\right)=g\left(\phi Z, \phi Z^{\prime}\right)+\sum \eta^{r}(Z) \eta^{r}\left(Z^{\prime}\right) ; Z, Z^{\prime} \in \Gamma T M$ and the fundamental 2 -form $\Omega$ associated with the $f$-structure satisfies:
(4) $\Omega\left(Z, Z^{\prime}\right)=g\left(\phi Z, Z^{\prime}\right) ; \Omega^{m} \wedge \varphi \neq 0, \varphi$ being the volume element of $M^{\perp}$, i.e., $\varphi=\eta^{2 m+1} \wedge \eta^{2 m+2} \wedge \eta^{2 m+3} \wedge \eta^{2 m+4}$.
Such a manifold $M\left(\phi, \Omega, \xi_{r}, \eta^{r}, g\right)$ is, as known, defined as framed $f$ manifold.

With respect to the cobasis $\mathcal{O}^{*}=\operatorname{covect}\left\{\omega^{A}, \eta^{r}\right\}$ the form $\Omega$ is expressed by $\Omega=\sum \omega^{a} \wedge \omega^{a^{*}} ; a \in\{1, \ldots, m\} ; a^{*}=a+m$ and the horizontal connection forms $\vartheta_{B}^{A}$ satisfies the known Kählerian conditions

$$
\begin{equation*}
\vartheta_{b}^{a}=\vartheta_{b^{*}}^{a^{*}} ; \vartheta_{b}^{a^{*}}=\vartheta_{a}^{b^{*}} \tag{5.1}
\end{equation*}
$$

Since on the other hand by (3.1) it is seen that the transversal connection forms $\vartheta_{A}^{r}$ vanish, one gets by exterior differentiation $d \Omega=0$. Since $\Omega$ is of constant rank and closed it follows that it is a presymplectic form on $M$ and a symplectic form on $M^{\top}$. We notice that in this case $\operatorname{ker}(\Omega)$ coincides with the vertical distribution $D_{p}^{\perp}$ of $M$ which may be also called characteristic distribution of $\Omega$. In addition by condition (3) of a framed $f$-structure and $\vartheta_{A}^{r}=0$ one has $(\nabla \phi) Z=0, Z \in \Gamma T M$, that is $\nabla$ and $\phi$ commute.

Recall now that the torsion tensor field $S$ of an f-structure is the vector valued 2 -form defined by $S=N_{\phi}+S^{\perp}$ where $N_{\phi}\left(Z, Z^{\prime}\right)=\left[\phi Z, \phi Z^{\prime}\right]+$ $\phi^{2}\left[Z, Z^{\prime}\right]-\phi\left[Z, \phi Z^{\prime}\right]-\phi\left[\phi Z, Z^{\prime}\right]$ is the Nijenhuis tensor field, and $S^{\perp}=$ $2 \sum d \eta^{r} \otimes \xi_{r}$ is the vertical component of $S$. By (3.10), (5.6) and $(\nabla \phi) Z=0$ it is easily seen that $S$ vanishes on $D^{\top}$. In this case, the $f$-structure $\left(\phi, \xi_{r}, \eta^{r}\right)$ is said to be horizontal-normal (or $D^{\top}$-normal) [2].

Consequently, following a definition of A. BEJANCU [2] the framed $f$ manifold $M\left(\phi, \Omega, \xi_{r}, \eta^{r}, g\right)$ under consideration is a framed- $C R$ manifold. On the other hand, taking into account that $\Omega$ is closed, the horizontal submanifold $M^{\top}$ of $M$ moves to a symplectic submanifold.

It also should be noticed that by (3.2) one may write $S^{\perp}$ as $S^{\perp}=$ $2 \omega \wedge I^{\perp} \Rightarrow d^{\nabla} S^{\perp}=0$ that is, $S^{\perp}$ is a closed vector valued 2 -form. We agree with the following

Definition 5.1. Let $M$ be a framed $f$-manifold and let $S^{\perp}$ be the vertical component of its associated torsion tensor. If the covariant differential of $S^{\perp}$ is closed, i.e., $d^{\nabla} S^{\perp}=0$, we say that $M$ is a vertical closed framed $f$-manifold.

Now since one finds $\mathcal{L}_{\mathcal{T}} \xi_{r}=\left[\mathcal{T}, \xi_{r}\right]=t_{r} \mathcal{T}-(t+a) \xi_{r}$ then one get at once $\mathcal{L}_{\mathcal{T}} S^{\perp}=2 \lambda S^{\perp}$. Accordingly the Lee vector field $\mathcal{T}$ defines an infinitesimal conformal transformation of $S^{\perp}$.

Then we can state the following
Theorem 5.2. Let $M\left(\phi, \Omega, \xi_{r}, \eta^{r}, g\right)$ be a framed $f$-manifold endowed with a $\mathcal{T}$-parallel almost contact 4-structure, and let $S^{\perp}$ be the vertical component of the torsion tensor field $S$ associated with the $f$-structure defined by $\phi$.

Any such $M$ is a framed $f$-CR manifold which is vertical torsion closed, i.e., $d^{\nabla} S^{\perp}=0$, and may be viewed as the local Riemannian product $M=M^{\top} \times M^{\perp}$ such that:
(i) $M^{\top}$ is a totally geodesic Kählerian submanifold of $M$, tangent to $\left\{\xi_{r}\right\}^{\perp}$;
(ii) $M^{\perp}$ is a totally geodesic space form submanifold of $M$, tangent to $\left\{\xi_{r}\right\}$;
(iii) the Lee vector field $\mathcal{T}$ of the ( $\mathcal{T} . P . A . C)$.4 -structure defines an infinitesimal conformal transformation of $S^{\perp}$.

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